Eikonal Theory

→ Rays, Eikonal Theory and Wave Propagation

eikonal (Greek) → icon → image

OV:

→ here, seek to provide description of wave propagation in 'short wavelength' limit [N.B. How short? – see HW on parabolic wave equation].

→ relevant to semi-classical limit of QM

→ description is in terms of rays – paths followed by wave

→ much of mechanics motivated by ray theory

Now:

– from previous HW, Fermat's minimum time principle (1662)

i.e.
$$T = \int_1^2 \frac{ds}{c(\underline{x})} = \frac{1}{c_0}\int_1^2 ds \; n(\underline{x})$$

T: travel time

n: index

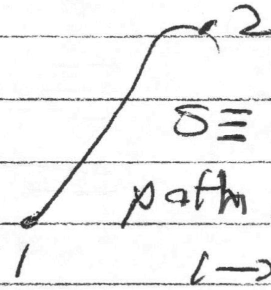
'ray Lagrangian'

$\delta T = 0 \Rightarrow$ ray path.

Generalizing the HW:

Fermat $\Rightarrow$

$$0 = \delta \int_1^2 n(\underline{x}(s))\, ds$$



$s \equiv$ ray path parameter $1 \to 2$ (dummy time)

$$= \delta \int_1^2 n(\underline{x}(s)) \left(\frac{d\underline{x}}{ds} \cdot \frac{d\underline{x}}{ds} \right)^{1/2} ds$$

$\equiv$ [redacted] $\int_1^2 L\, ds$

$\Rightarrow$

$$0 = \int_1^2 \left(\frac{\partial L}{\partial \underline{x}} \cdot \delta\underline{x} + \frac{\partial L}{\partial \left(\frac{d\underline{x}}{ds}\right)} \cdot \delta\left(\frac{d\underline{x}}{ds}\right) \right)$$

$$= \underset{\text{chp}}{\text{e.p.}} + \int_1^2 \left(\frac{\partial L}{\partial \underline{x}} \cdot \delta\underline{x} - \frac{d}{ds}\left(\frac{\partial L}{\partial\left(\frac{d\underline{x}}{ds}\right)} \right) \cancel{\delta\underline{x}} \right)$$

$\Rightarrow$

$$\boxed{\begin{array}{l} \dfrac{\partial L}{\partial \underline{x}} - \dfrac{d}{ds}\left(\dfrac{\partial L}{\partial\left(\frac{d\underline{x}}{ds}\right)} \right) = 0 \\[2ex] L = n(\underline{x}(s)) \left(\dfrac{d\underline{x}}{ds} \cdot \dfrac{d\underline{x}}{ds} \right)^{1/2} \end{array}}$$

Crank $\Rightarrow$

if $|\dot{\underline{x}}| = \left[\frac{d\underline{x}}{ds} \cdot \frac{d\underline{x}}{ds} \right]^{1/2}$

$$\boxed{ |\dot{\underline{x}}| \frac{\partial n}{\partial \underline{x}} - \frac{d}{ds} \left(n(\underline{x}) \frac{\dot{\underline{x}}}{|\dot{\underline{x}}|} \right) = 0 }$$

→ general expression

→ $\partial n / \partial \underline{x}$ ⟺ effective force on ray
($U \Leftrightarrow n$)

→ $n(\underline{x}) \frac{\dot{\underline{x}}}{|\dot{\underline{x}}|}$ ⟺ defines generalized momentum analogue.

$\left(n(\underline{x}) \frac{d\underline{x}}{ds} \right)$

Note: $ds^2 = d\underline{x} \cdot d\underline{x}$
so $|\dot{x}| = 1$

$$\Rightarrow \boxed{ \frac{\partial n}{\partial \underline{x}} - \frac{d}{ds} \left(n(\underline{x}) \frac{d\underline{x}}{ds} \right) = 0 }$$

is equivalent.

→ A bit of geometry:

$$\frac{d}{ds}\left(n(\underline{x})\,\frac{d\underline{x}}{ds}\right) - \frac{\partial n}{\partial \underline{x}} = 0$$ → ray equation

⇒

$$n(\underline{x})\frac{d^2\underline{x}}{ds^2} + \left(\frac{\partial n}{\partial \underline{x}} \cdot \frac{d\underline{x}}{ds}\right)\frac{d\underline{x}}{ds} = \frac{\partial n}{\partial \underline{x}}(x)$$

∴

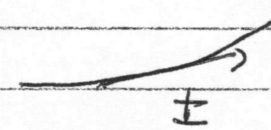
$$\boxed{\frac{d^2\underline{x}}{ds^2} = \frac{1}{n(\underline{x})}\frac{\partial n}{\partial \underline{x}} - \frac{1}{n(\underline{x})}\left(\frac{\partial n}{\partial \underline{x}} \cdot \frac{d\underline{x}}{ds}\right)\underline{t}} \qquad t = \frac{dx}{ds}$$

What does it mean?

→ $d\underline{x}/ds$ is unit tangent to ray.

i.e. $ds\,ds = d\underline{x} \cdot d\underline{x}$

$\underline{t} = d\underline{x}/ds$



so

⇒ $d^2\underline{x}/ds^2$ corresponds to ray curvature κ.

$1/|\underline{K}| \equiv$ effective radius of curvature

$\underline{K} \equiv$ curvature

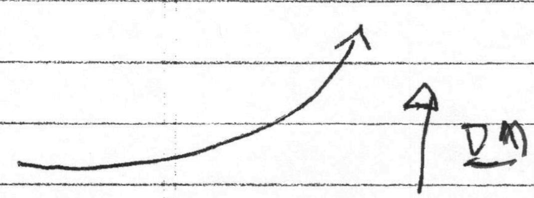
so

$$\underline{K} = \frac{1}{n}\underline{\nabla}n - \frac{1}{n}(\underline{t}\cdot\underline{\nabla}n)\,\underline{t}$$

$$= \frac{1}{n}\left(\underline{\nabla}n \cdot \underline{\hat{n}}_0\right)\underline{\hat{n}}_0$$

unit normal to path

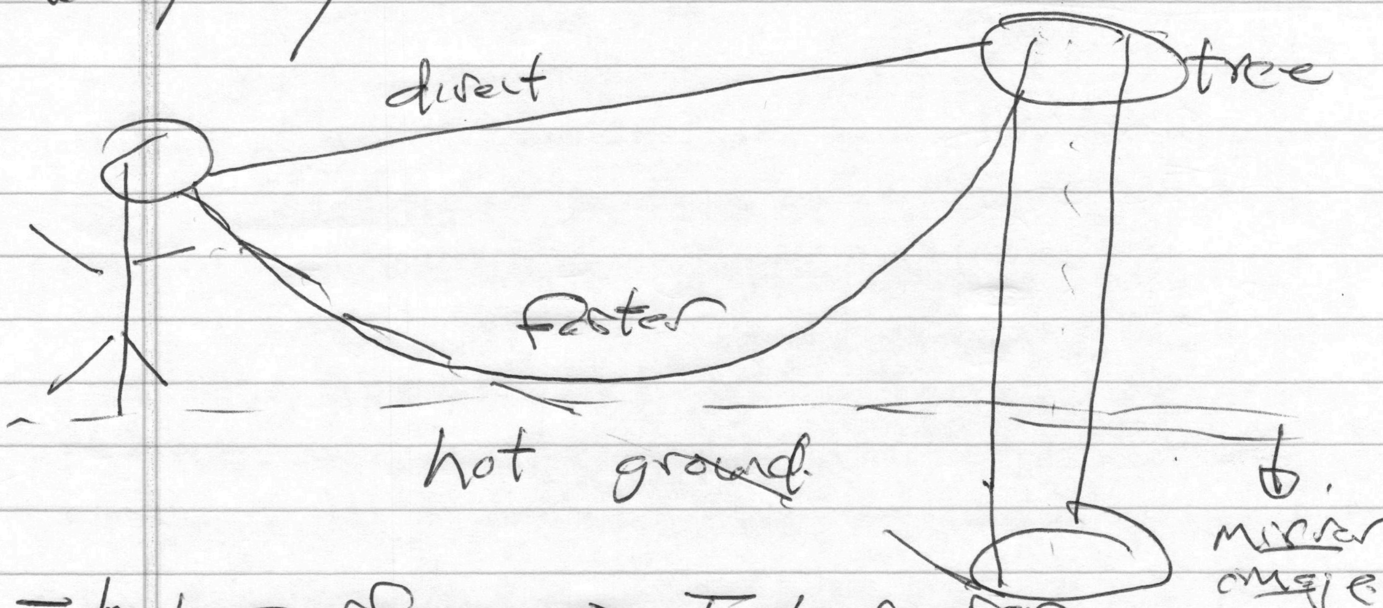
Loosely put, ray curves toward region of increasing index.



→ Mirages

- mirages are optical illusions of reflection from water, etc. which occur outdoors.

- how/why



- hot surface ⇒ T decreases
  air density increases with height

- index n ~ density

- so, observer sees direct path image and curved path-induced mirror image

- if no tree → blue sky → appears like water → mirage

3c': [redacted]

- so reasonable to take index ~ z

$$n(z) = n_0 (1 + \alpha z)$$

Now, Fermat $\Rightarrow$ ray from:

$$\delta \int \left(1 + (dz/dx)^2\right)^{1/2} n(z) = 0$$

3d.

So

$$\frac{d}{dx}\left(\frac{n(z)}{\left(1+(dz/dx)^2\right)^{1/2}}\frac{dz}{dx}\right) = \left(1+\left(\frac{dz}{dx}\right)^2\right)^{1/2}\frac{dn}{dz}$$

$dz/dx \equiv \dot{z}$

$\Rightarrow$

$$\frac{d}{dx}\left(\frac{n_0(1+\alpha z)}{(1+\dot{z}^2)^{1/2}}\dot{z}\right) = n_0(1+\dot{z}^2)^{1/2}\alpha$$

For ~~[illegible]~~ $\approx$ horizontal rays,

$\dot{z}^2 \ll 1$

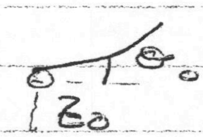
$\alpha z \ll 1$

$\Rightarrow$

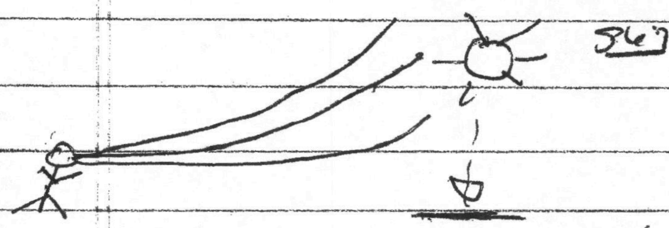
$$\frac{d^2 z}{dx^2} \cong \alpha$$

$\therefore$ then have:

$$z(x) = \left(\frac{\alpha}{2}x^2 + \underbrace{\tan\theta_0}_{\text{inclination}}\, x + z_0\right)$$



then rays diverge parabolically,



apparent location
(shimmering, bright light)

⇒ mirage
(appears like reflection
from water)

Origin of shimmer? ⟹ conv. turbulence.

Now, consider:

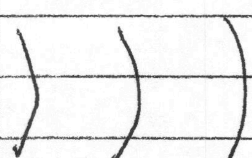
→ Helmholtz Eqn.

$$\nabla^2\psi + \frac{\omega^2}{c(x)^2}\psi = 0$$

→ index

$$1/c(x)^2 \equiv \frac{n(x)^2}{c_0^2}$$ → ref. speed

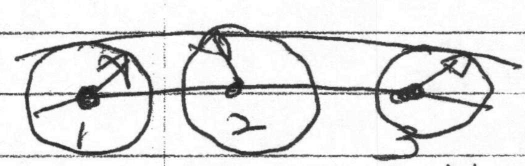
→ consider phase front



ϕ = const. surfaces

↓ phase

Now, to describe propagation:



ϕ = const. surf. at $t+\Delta t$

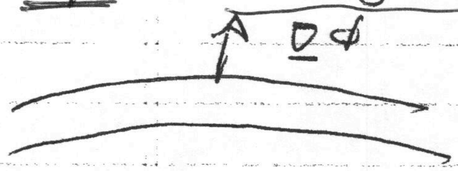
ϕ = const surface at t_0

$c_1\Delta t$ $c_2\Delta t$ $c_3\Delta t$

i.e. each point on surface ϕ = const at t emits spherical disturbance.

sum of spherical disturbances defines new constant phase surface. Curvature due $c(\underline{x})$.
Envelope of spheres $\Rightarrow$ wavefront at $t_0 + \Delta t$

- rays orthogonal to wave fronts.



$\Rightarrow$ root motivation of Hamiltonian Mech.

Now, $\forall$ infinitesimal displacement vector along ray $\equiv d\underline{\sigma}$

i.e. $d\underline{\sigma} \parallel \underline{\nabla}\phi$

then, since equivalent to advance in space or time,

$$\underline{\nabla}\phi \cdot d\underline{\sigma} = \omega \, dt$$

$$|\underline{\nabla}\phi| \, |d\sigma| = \omega \, dt$$

$dt = d\sigma / c$ (by definition)

$$\Rightarrow \quad |\underline{\nabla}\phi| \, |d\sigma| = \omega \frac{d\sigma}{c}$$

$\therefore$

$$|\underline{\nabla}\phi| = \omega/c$$

$$\Rightarrow \boxed{(\underline{\nabla}\phi)^2 = \omega^2/c^2}$$

— eikonal equation

$\Rightarrow$ eqn. for spatial evolution ϕ.

[reduces wave eqn to phase eqn.

N.B. — Can obtain directly from Helmholtz Eqn.

$$\nabla^2\psi + \frac{\omega^2}{c(x)^2}\psi = 0$$

$$\psi = A\, e^{i\phi(x)/\epsilon} \quad (\text{WKB})$$

$\epsilon \to 0$ (short wavelength)

$\Rightarrow$

$$\left[-\frac{(\underline{\nabla}\phi)^2 A}{\epsilon^2} + i\frac{\nabla^2\phi\, A}{\epsilon} + 2i\frac{\underline{\nabla}A\cdot\underline{\nabla}\phi}{\epsilon} + \nabla^2 A \right] e^{i\phi} = -\omega^2/c(x)^2\, A\, e^{i\phi}$$

so dominant balance i

$$+\frac{(\underline{\nabla}\phi)^2}{\epsilon^2} = \frac{\omega^2}{c(\underline{x})^2}$$

- now absorb ϵ to ϕ.

- note eikonal lowers order of problem $\Rightarrow$ first order pde.

Now, by construction

$$\underline{\nabla}\phi \cdot d\underline{l} \equiv \text{net phase increment along ray.}$$

so $$\underline{\nabla}\phi = \underline{k} = \underline{k}(\underline{x})$$

in sense of WKB.

(n.b. generally, $\partial\phi/\partial t = -\omega$).

$$\phi = \int \underline{k}\cdot d\underline{x} = \int \underline{\nabla}\phi \cdot d\underline{x}$$

$$= \int \underline{k}\cdot d\underline{l}$$

$$\psi = A \exp\left[i\left(\int \underline{k}\cdot d\underline{x} - \omega t\right)\right]$$

is eikonal approximation to wave fctn.

- N.B. $\Rightarrow$ $\underline{k}$ specifies ray direction orthog. to ϕ (phase) surfaces.

$\Rightarrow$ Now, seek equations which evolve ray path in time, space i.e. give
- ray position $\underline{x}$ as fctn of
- ray direction $\underline{k}$

time.

$\Rightarrow$ defines mechanics problem.

a.) Poor Man's Version

- For linear wave, have $\omega = \text{const.}$

Since $\omega = \omega(\underline{k}, \underline{x}) \Rightarrow \overbrace{(\underline{x}, \underline{k})}^{\text{ray}}$

$$\frac{d\omega}{dt} = 0 = \cancel{\frac{\partial\omega}{\partial t}} + \frac{\partial\omega}{\partial\underline{k}}\cdot\frac{d\underline{k}}{dt} + \frac{\partial\omega}{\partial\underline{x}}\cdot\frac{d\underline{x}}{dt}$$

$$\Rightarrow \boxed{\begin{aligned} d\underline{k}/dt &= -\partial\omega/\partial\underline{x} \\ d\underline{x}/dt &= \partial\omega/\partial\underline{k} = \underline{V}_{gr} \end{aligned}}$$

eikonal equations

Hamiltonian EOMs

With, of course:

$$\omega^2 = c(\underline{x})^2 k^2$$

$$\cancel{2}\omega\,\partial\omega = \cancel{2}\underline{k}\cdot\partial\underline{k}\;\; c(\underline{x})^2$$

$$\underline{k} = k\hat{k}$$

$$\hat{k} = \frac{\nabla\phi}{|\nabla\phi|}$$

$$\partial\omega = \hat{k}\cdot\partial\underline{k}\;\; c(x)$$

$$\partial\omega/\partial\underline{k} = c(\underline{x})\,\hat{k}$$

$\equiv$ group velocity.

$$\frac{\partial\omega}{\partial\underline{x}} = \frac{\partial}{\partial\underline{x}}\left[c(x)^2 k^2\right]^{1/2} = k\,\frac{\partial c(\underline{x})}{\partial\underline{x}}$$

$\therefore$

$$\boxed{\begin{aligned} \frac{d\underline{x}}{dt} &= c(\underline{x})\,\hat{k} \\ \frac{d\underline{k}}{dt} &= -k\,\frac{\partial c(\underline{x})}{\partial\underline{x}} \end{aligned}}$$

$c(\underline{x})$ profile determines ray path.

eikonal equation for acoustics.

b.) More Rigorously ----

$$\Phi = \int\left[\underline{k}\cdot d\underline{x} - \omega\,dt\right] \rightarrow \text{total phase}$$

$$dS = L\,dt$$
$$= (p\dot{q} - H)\,dt$$

$$d\Phi = \underline{k}\cdot d\underline{x} - \omega\,dt = (\underline{k}\cdot\dot{\underline{x}} - \omega)\,dt$$

Now, assert ray will follow path which extremizes Φ, i.e. {minimizes accumulated phase?}.

Note: analogy of phase and action

: later demonstrate connection to Fermat

$$\delta\Phi = \delta\int\left[\underline{k}\cdot d\underline{x} - \omega\,dt\right] = 0$$

$$= \int\left\{\left[\delta\underline{k}\cdot d\underline{x} + \underline{k}\cdot\delta d\underline{x}\right] - \left(\frac{\partial\omega}{\partial\underline{k}}\cdot\delta\underline{k} + \frac{\partial\omega}{\partial\underline{x}}\cdot\delta\underline{x}\right)dt\right\}$$

as usual, $\delta\underline{x} = \delta\underline{k} = 0$, at end points.

So integrating by parts:

$$\delta\Phi = \int\left\{\left[\delta\underline{k}\cdot d\underline{x} - d\underline{k}\cdot\delta\underline{x}\right] + \text{e.p.} - \left[\left(\frac{\partial\omega}{\partial\underline{k}}\cdot\delta\underline{k}\right) + \left(\frac{\partial\omega}{\partial\underline{x}}\cdot\delta\underline{x}\right)\right]dt\right\}$$

so $d\underline{x} = \left(\frac{\partial\omega}{\partial\underline{k}}\right) dt$

$$d\underline{k} = -\left(\frac{\partial\omega}{\partial\underline{x}}\right) dt$$

$$\Rightarrow \boxed{\begin{aligned} \frac{d\underline{x}}{dt} &= \frac{\partial\omega}{\partial\underline{k}} \\ \frac{d\underline{k}}{dt} &= -\frac{\partial\omega}{\partial\underline{x}} \end{aligned}}$$

→ eikonal equations

Note:

→ evolve ray in $\underline{x}$, $\underline{k}$ phase space.

$\underline{x}$ → position; $\underline{k}$ → direction, momentum

→ <u>Hamiltonian</u> <u>equations</u> for ray in $(\underline{x}, \underline{k})$ phase space.

i.e. $\frac{\partial}{\partial\underline{x}} \cdot \frac{d\underline{x}}{dt} + \frac{\partial}{\partial\underline{k}} \cdot \frac{d\underline{k}}{dt}$

$$= \frac{\partial}{\partial\underline{x}} \cdot \frac{\partial\omega}{\partial\underline{k}} - \frac{\partial}{\partial\underline{k}} \cdot \frac{\partial\omega}{\partial\underline{x}} = 0.$$

→ since eikonal equations Hamiltonian, can define:

$\rho(\underline{x}, \underline{k}, t) \equiv$ wave density in $\underline{x}, \underline{k}$ phase space.

$N(\underline{x}, \underline{k}, t)$
- wave action density
- ~ Wigner dist.
- ~ intensity.

and use Liouville's Thm:

$k \to$ ladder

Pull in

$$\boxed{\frac{\partial \rho}{\partial t} + \underline{v}_{gr} \cdot \frac{\partial \rho}{\partial \underline{x}} - \frac{\partial \omega}{\partial \underline{x}} \cdot \frac{\partial \rho}{\partial \underline{k}} = 0}$$

$$\frac{\partial \rho}{\partial t}(\underline{x}, t) + \nabla \cdot [\underline{v}_{gr} \rho] = 0$$

$\partial_t \omega = 0$

$\partial_t \mathcal{E} + \nabla \cdot (\underline{v}_{gr} \mathcal{E}) = 0$

- wave kinetic eqn.
- relates ρ, and intensity, to $C(\underline{x})$ profiles, for acoustics
- gives intensity evol'n.
- applications in radiation hydro, quasi-particle evolution, etc.

Obvious analogy: (Hamiltonian Systems)

Particles	Rays
H	ω
$\underline{p}$	$\underline{k}$
$\underline{q}$	$\underline{x}$
S	ϕ

→ Eikonal Theory; Supplement

Recall:

- for Helmholtz Eqn, derived:

$$(\nabla\phi)^2 = \frac{\omega^2}{c(x)^2} \quad \Rightarrow \text{ eikonal eqn.}$$

↳ inhomogeneous speed.

$\psi \sim A e^{i\phi}$

- as rays ⊥ phase fronts

$$\nabla\phi \equiv \underline{k} = k(x)$$

aka WKB

$$\omega = -\frac{\partial\phi}{\partial t}$$

so

$$\psi = A \exp\Big[\underbrace{i\Big(\int \underline{k}(\underline{x})\cdot d\underline{x} - \omega t\Big)}_{\Phi} \Big]$$

eikonal approx. to wave function

- now, for ray trajectories, observe total phase Φ

$$d\Phi = \underline{k}\cdot d\underline{x} - \omega\, dt$$

$$= \left(\underline{k}\cdot\frac{d\underline{x}}{dt} - \omega\right) dt$$

$S \leftrightarrow \Phi$ is key analogy

analogous

$$S = \int L\, dt \Rightarrow dS = L\,dt$$
$$= (p\dot{q} - H)\, dt$$

$\therefore$

— obvious analogy

$\underline{k} \leftrightarrow \underline{p}$

$\underline{x} \leftrightarrow \underline{q}$

$\omega \leftrightarrow H$

i.e. QM: $\underline{p} = \hbar \underline{k}$

$E = \hbar\omega$

so $\dfrac{d\underline{k}}{dt} = -\dfrac{\partial\omega}{\partial\underline{x}} \Leftrightarrow \dfrac{d\underline{p}}{dt} = -\dfrac{\partial H}{\partial\underline{q}}$

$\frac{dp}{dt} = -\partial H/\partial q$

$\dfrac{d\underline{x}}{dt} = \dfrac{\partial\omega}{\partial\underline{k}} \Leftrightarrow \dfrac{d\underline{q}}{dt} = \dfrac{\partial H}{\partial\underline{p}}$

$\frac{dq}{dt} = \frac{\partial H}{\partial p}$

eg in terms $c(\underline{x})$:

$$\omega^2 = c(x)^2 k^2$$

$$\frac{d\underline{k}}{dt} = -k \frac{\partial \underline{c}(\underline{x})}{\partial \underline{x}}$$

$$\frac{d\underline{x}}{dt} = c(\underline{x})\, \hat{k} .$$

N.B.

→ $\partial\omega/\partial\underline{k} \equiv \underline{v}_{gr}$ group velocity

What does $\underline{v}_{gr}$ mean ?

Consider wave packet, — carrier $\underline{k}_0$ spread $\Delta \underline{k}$

$$\phi|_{t=0} \sim \underbrace{e^{i\underline{k}_0\cdot\underline{x}}}_{\text{carrier}}\, \underbrace{F(\underline{x})}_{\text{envelope}}$$

$\underline{k}_0$ → carrier

$$F(\underline{x}) \sim \sum_{\Delta \underline{k}} \underbrace{e^{i\Delta\underline{k}\cdot\underline{x}}}_{\text{envelope}} .$$

so

$$\phi(\underline{x},t) \sim \sum_{\Delta\underline{k}} e^{i[(\underline{k}_0+\Delta\underline{k})\cdot\underline{x} - \omega(\underline{k}_0+\Delta\underline{k})t]}$$

$$\sim \underbrace{e^{i(\underline{k}_0\cdot\underline{x} - \omega(\underline{k}_0)t)}}_{\text{carrier}} \sum_{\Delta\underline{k}} e^{i\Delta\underline{k}\cdot\underline{x}}\, e^{-\frac{\partial\omega}{\partial\underline{k}}\cdot\Delta\underline{k}\, t}$$

$$\phi(\underline{x},t) \sim e^{i(\underline{k}_0 \cdot \underline{x} - \omega t)} \, F\Big(\underline{x} - \underbrace{\frac{\partial \omega}{\partial \underline{k}}}_{\underline{v}_{gr}} t\Big)$$

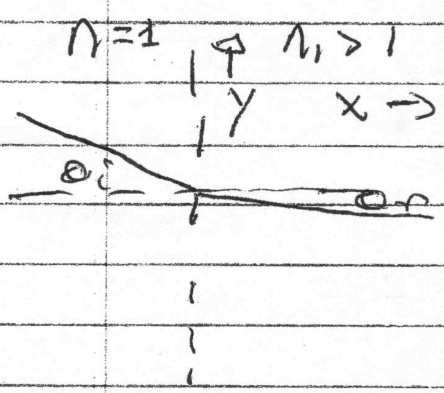
$\Rightarrow$ rate/speed at which energy propagated. $\rightarrow |\phi|^2$

N.B. $\underset{pckt}{E} \sim |\phi_{\underline{k}}|^2 \sim |f|^2$

$\Rightarrow$ V_{gr} sets speed at which energy propagates.

$\rightarrow \dfrac{d\underline{k}}{dt} = -\dfrac{\partial \omega}{\partial \underline{x}}$ $\Rightarrow$ Snell's Law.

i.e.



$$\frac{d\underline{k}}{dt} = -\frac{\partial \omega}{\partial \underline{x}} \quad \Rightarrow \quad \frac{dk_y}{dt} = 0$$

$$k_{y_-} = k_{y_+} \quad \Rightarrow \quad k_- \sin\theta_i = k_+ \sin\theta_r$$

$$k_-^2 = n_0^2 \frac{\omega^2}{c_0^2} \qquad k_+^2 = n_1^2 \omega^2 / c_0^2$$

$n_0 \sin\theta_i = n_i \sin\theta_r$ ✓.

- Now, if seek first principles approach

$\Rightarrow$ extremize Φ (i.e. look for phase stationarity)

stationarity $\to$ trajectories i.e. ray or particle (cf. full wave)

$$\delta\Phi = \delta\int[\underline{k}\cdot d\underline{x} - \omega dt]$$

$$= \delta\int[\underline{k}\cdot\dot{\underline{x}} - \omega]\,dt$$

($t \to iT$ steepest descents; $\int d\underline{k}\, e^{i\Phi} \to$ packet)

$$= \int\left[\delta\underline{k}\cdot\dot{\underline{x}} + \underline{k}\cdot\delta\dot{\underline{x}} - \frac{\partial\omega}{\partial\underline{x}}\cdot\delta\underline{x} - \frac{\partial\omega}{\partial\underline{k}}\cdot\delta\underline{k}\right]dt$$

but $\delta\dot{\underline{x}} = \frac{d}{dt}\delta\underline{x}$

e.p. fixed.

$$\delta\Phi = \cancel{\underline{k}\cdot\delta\underline{x}\Big|_{t_1}^{t_2}} + \int\Big[\delta\underline{k}\cdot\dot{\underline{x}} - \frac{d\underline{k}}{dt}\cdot\delta\underline{x}$$

$$- \frac{\partial\omega}{\partial\underline{x}}\cdot\delta\underline{x} - \frac{\partial\omega}{\partial\underline{k}}\cdot\delta\underline{k}\Big]$$

$\delta \Phi \Rightarrow$

$$\frac{dx}{dt} = \frac{\partial \omega}{\partial k}, \qquad \frac{dk}{dt} = -\frac{\partial \omega}{\partial x}.$$

$\Rightarrow$ Liouville Thm $\rightarrow$ Wave Kinetics.

N.B.: for semi-classical limit

$$P = N \hbar \underline{k}$$

$$N \Leftrightarrow \rho.$$

$$E = N \hbar \omega$$

Finally, : to recover Fermat, note:

$$\delta \Phi = 0$$

$$d\phi = \underline{k} \cdot d\underline{x} - \omega t$$

so for ray path:

$$\delta \Phi_s = 0 \, , \qquad \Phi_s = \int \underline{k} \cdot d\underline{x}$$

but

$$\underline{k} \cdot d\underline{x} = \underline{k} \cdot \frac{d\underline{x}}{ds} \, ds \Rightarrow$$

$$\underline{k} = \underline{\nabla} \phi = |\underline{\nabla} \phi| \, \underline{t}$$

$$d\underline{x}/ds = \underline{t}$$

so

$$\delta \Phi_s = \delta \int |\nabla \phi| \, ds$$

but $|\nabla \phi|^2 = \omega^2 / c^2 = \frac{\omega^2}{c_0^2} n(\underline{x})^2$

$\Rightarrow$

$$\delta \Phi_s = \delta \frac{\omega}{c_0} \int ds \, n(\underline{x}) = 0$$

$$\Rightarrow \delta \int ds \, n(\underline{x}) = 0$$